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Published in:
Advanced Energy Materials

DOI:
[10.1002/aenm.201100252](https://doi.org/10.1002/aenm.201100252)

IMPORTANT NOTE: You are advised to consult the publisher's version (publisher's PDF) if you wish to cite from it. Please check the document version below.

Document Version
Publisher's PDF, also known as Version of record

Publication date:
2011

[Link to publication in University of Groningen/UMCG research database](#)

Citation for published version (APA):

Dung, N. H., Ou, Z. Q., Caron, L., Zhang, L., Cam Thanh, D. T., Wijs, G. A. D., Groot, R. A. D., Buschow, K. H. J., & Brück, E. (2011). Mixed Magnetism for Refrigeration and Energy Conversion. *Advanced Energy Materials*, 1(6), 1215-1219. <https://doi.org/10.1002/aenm.201100252>

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Mixed Magnetism for Refrigeration and Energy Conversion

Nguyen H. Dung, Zhi Qiang Ou, Luana Caron, Lian Zhang, Dinh T. Cam Thanh, Gilles A. de Wijs, Rob A. de Groot, K. H. Jürgen Buschow, and Ekkes Brück*

The efficient coupling between lattice degrees of freedom and spin degrees of freedom in magnetic materials can be used for refrigeration and energy conversion. This coupling is enhanced in materials exhibiting the giant magnetocaloric effect. First principle electronic structure calculations on hexagonal MnFe(P, Si) reveal a new form of magnetism: the coexistence of strong and weak magnetism in alternate atomic layers. The weak magnetism of Fe layers (disappearance of local magnetic moments at the Curie temperature) is responsible for a strong coupling with the crystal lattice while the strong magnetism in adjacent Mn-layers ensures Curie temperatures high enough to enable operation at and above room temperature. Varying the composition on these magnetic sublattices gives a handle to tune the working temperature and to achieve a strong reduction of the undesired thermal hysteresis. In this way we design novel materials based on abundantly available elements with properties matched to the requirements of an efficient refrigeration or energy-conversion cycle.

a thermomagnetic generator.^[3] This device can be used to generate electricity from 'waste heat', heat that currently is released unused into the environment. Efficient magnetocaloric materials could therefore contribute significantly to a reduction in energy consumption. However, up to now no suitable magnetic materials were available to operate between room temperature and 400K.

Following the discovery of giant MCE in $\text{Gd}_5(\text{Si,Ge})_4$,^[4] a number of magnetocaloric materials^[4–10] with a first-order magnetic phase transition (FOMT) have been intensively explored. In these materials, the FOMT enhances the magnetocaloric effect in the vicinity of the magnetic phase transition. The maximum isothermal entropy change is therefore often significantly greater than that of the benchmark material $\text{Gd}^{[11]}$ that presents a second-order

1. Introduction

Limited resources and the wish for improved prosperity call for efficient use of energy. The UN Advisory Group on Energy and Climate Change recommends a target of 40% improved efficiency by 2030.^[1] Materials research can contribute significantly to reach this target. Magnetic refrigeration offers potential to achieve a 50% higher energy-efficiency compared to vapour compression refrigeration.^[2] This makes magnetic refrigeration a technology that attracts growing attention. Magnetic refrigeration is based on the magnetocaloric effect (MCE); solid magnetic materials heat up or cool down when an external magnetic field is applied or removed, due to the entropy transfer ΔS between the magnetic system and the crystal lattice. In a reversed process heat can be transformed into electricity with

magnetic phase transition. Combining giant magnetocaloric materials with different T_C in series, a higher efficiency and a greater temperature span than that of Gd is obtained.^[12] For optimal performance the materials used in such a composite regenerator need to have very similar magnetocaloric properties, to achieve a constant entropy change as function of temperature.^[13] The large thermal hysteresis frequently associated with the FOMT seriously hampers the application in a refrigeration cycle.^[14] Thermal and field hystereses are intrinsic properties of a first-order material. However, the size of this hysteresis may strongly depend on microstructure or strain in the system. Here we report on a novel mechanism for FOMT derived from first principle electronic structure calculations. The intercalation of weak and strong magnetism in adjacent lattice planes induces a large magneto-elastic effect and giant magnetocaloric effects. With this mechanism we can generate exceptionally favourable magnetocaloric properties in broad regions of the Mn-Fe-P-Si system.

2. Results and Discussion

Fe_2P and related alloys have attracted attention for quite some time,^[15] on replacing more than 10% of P by Si in Fe_2P the hexagonal crystal lattice is transformed into an orthorhombic one.^[16] To avoid the hexagonal to orthorhombic transformation we replaced Fe on the 3g sites by Mn to form a compound with very interesting properties. $\text{MnFeP}_{0.5}\text{Si}_{0.5}$ with a FOMT near room temperature, crystallizes in the hexagonal Fe_2P type

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DOI: 10.1002/aenm.201100252

structure that has four distinct lattice sites, the thermal hysteresis ΔT_{hys} of 35 K makes this material however unsuitable for applications.^[17] In contrast to most other magnetocaloric materials the volume change in this material is rather small and we observe mainly a change in c/a ratio of the hexagonal lattice. In order to elucidate the origin of the observed magnetocaloric effect, electronic structure calculations were performed on the ferromagnetic ground-state, while the behaviour at the Curie temperature was modelled by a supercell obtained by doubling the unit cell (allowing for antiferromagnetic configurations). The calculations show that layers occupied by manganese are strongly magnetic; implying that the magnetic order only is lost at the Curie temperature. The size of the Mn moment is reduced from $2.8 \mu_B$ in the ferromagnetic phase to $2.6 \mu_B$ in the paramagnetic phase. By contrast, the iron-layers show weak itinerant magnetism: here the Fe moment in the ferromagnetic phase is $1.54 \mu_B$, while in the paramagnetic phase it vanishes ($\sim 0.003 \mu_B$). This implies that the electron density around the Fe sites changes drastically at the phase transition. This change is especially significant within the Fe layer as illustrated in **Figure 1**. In this figure we show the difference in electron density between the ferromagnetic and the paramagnetic state. The dominant changes occur close to the Fe sites while near the Si/P sites the electron density is hardly affected. In the ferromagnetic state high electron density forms a dumbbell pointing into the empty space between adjacent Si/P atoms, while in the paramagnetic state high electron density forms a clover four pointing towards the nearest Si/P neighbours. This redistribution of electron density means that non-bonding electron density at the Fe site below T_C changes into a distribution which

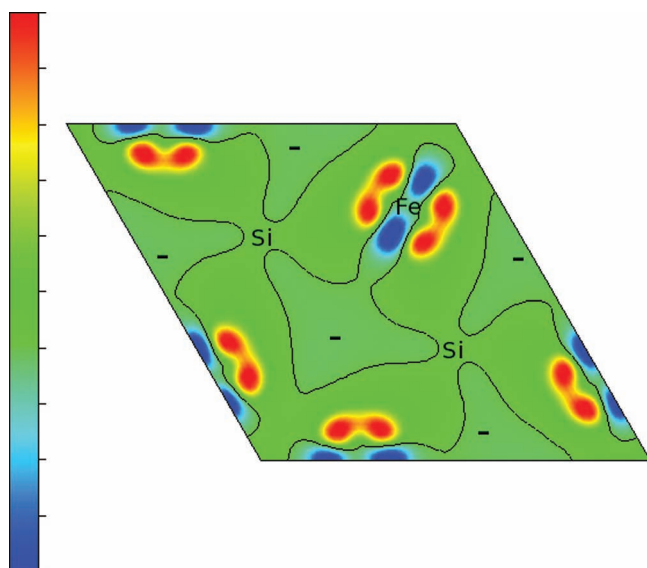


Figure 1. Difference between the electron densities calculated for situations above and below the Curie temperature in the iron-silicon/phosphorous plane (ferromagnetic density (–) subtracted from the paramagnetic one). This results locally in negative electron densities where the highest ferromagnetic density was. The colour codes for the electron density ranges from -4.55×10^{-5} (dark blue) to 9.1×10^{-5} electrons/ \AA^3 at the red end of the scale. The dark lines indicate no change in electron densities.

is hybridized with Si/P above T_C . This change in hybridization causes the distinct change in c/a ratio observed experimentally in the magneto-elastic transition at T_C .

The loss of moments on the iron site is also clear from the partial density of states as function of energy shown in **Figure 2**. It shows identical curves for the two spin directions for iron above the Curie temperature, in sharp contrast with manganese that maintains its moment.

Such a combination of strong and weak magnetism in *one and the same* compound is unexpected. Weak magnetism is rare; it is found in materials like ZrZn_2 ^[18] or Ni_3Al .^[19] Curie

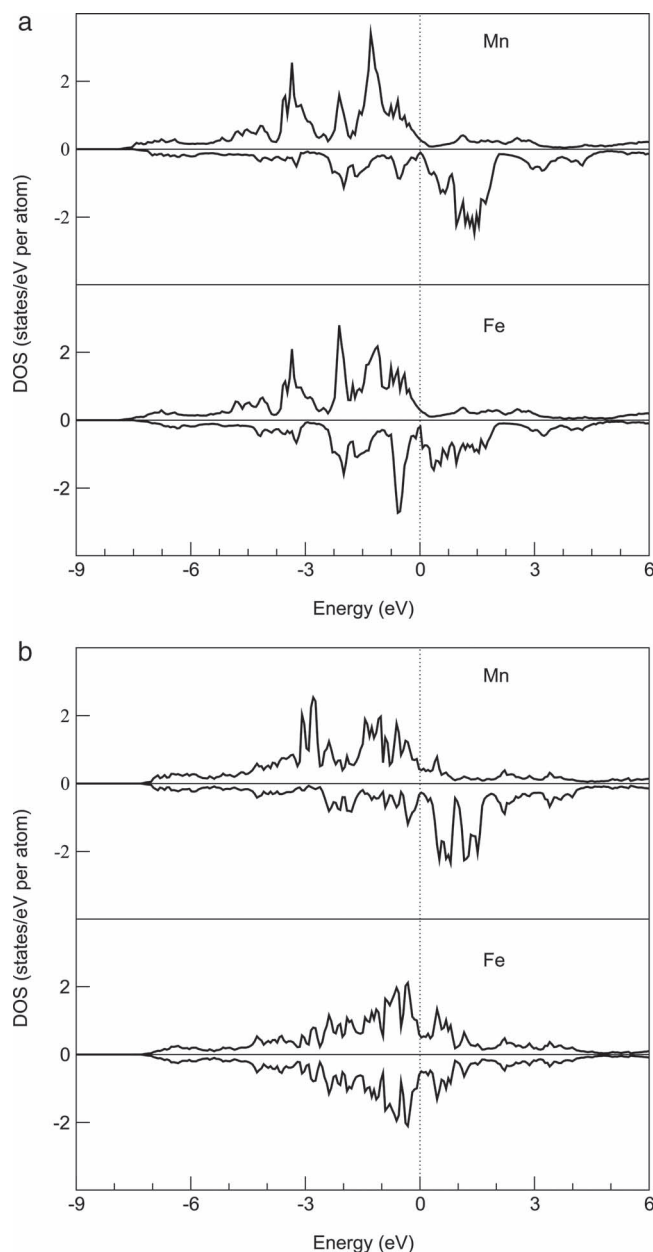


Figure 2. Partial local densities of states for the magnetic atoms in the ferromagnetic state (a) and the paramagnetic state (b) representing the state above the Curie point. Note the identical densities of state for the iron above the Curie temperature.

temperatures are low (for example, ZrZn_2 : 33 K; Ni_3Al : 23–58 K depending on composition). To the best of our knowledge no other cases of mixed magnetism have been described before. It is directly related to the giant magneto-caloric effect, because in solids the existence of magnetic moments competes with chemical bonding.^[20] The loss of the magnetic moments of iron enables the strong coupling to the lattice above the Curie temperature resulting in the discontinuity of the c/a ratio leading to the FOMT. On the other hand, the strong magnetism of the manganese layers ensures a Curie temperature near room temperature.

Experimentally we studied the effect of changing the lattice sites occupations. The hexagonal Fe_2P type of structure is found to be stable over a broad range of compositions. Similar to the orthorhombic compound^[16] T_C increases with increasing Si content. This increase in T_C is probably caused by an increase in magnetic moments that is seen from our theoretical calculations. These calculations indicate that mainly the Fe moment on 3f sites is enhanced. Neutron diffraction experiments are planned to test this finding experimentally. When increasing the Mn content in $\text{MnFe}(\text{P}, \text{Si})$ compounds, T_C decreases (see Figure 3a). This finding may be astonishing at the first glance as Mn appears to carry a larger moment than Fe. However, the excess Mn atoms have to occupy the 3f sites now where they can develop only a relatively low moment of $1.25 \mu_B$ compared with $1.54 \mu_B$ of Fe atoms on these sites. Additionally, it is well known that Mn alloys tend to order antiferromagnetically below a critical Mn–Mn distance. The distance between 3g and 3f sites is clearly below this critical distance and therefore Mn on 3f sites will not contribute to the strong ferromagnetism. On the other hand an increase of the Fe content results in an increase of T_C . The excess Fe will occupy the 3g sites where Fe always carries a large magnetic moment, and T_C of hexagonal $\text{Fe}_2\text{P}_{0.8}\text{Si}_{0.2}$ 510 K exceeds the T_C of the Mn containing alloy^[16]

Both Fe substitutions on the Mn sublattice or Mn substitutions on the Fe sublattice, as well as an increase in Si content, are beneficial in that they give rise to a decrease in ΔT_{hys} (Figure 3b). From these trends we derive that a large ΔS_m coupled with a small ΔT_{hys} can be obtained by balancing the Mn:Fe ratio and the P:Si ratio. Furthermore, T_C can be tuned by changing the Mn:Fe and P:Si ratios simultaneously to keep both a large ΔS_m and a small ΔT_{hys} . These trends also hold for slightly non-stoichiometric compounds. By concurrently changing Mn:Fe and P:Si ratios in $\text{Mn}_x\text{Fe}_{1.95-x}\text{P}_{1-y}\text{Si}_y$ compounds, the working temperature can be controlled between 210 and 430 K for $x = 1.35$, $y = 0.46$ and $x = 0.66$ and $y = 0.42$, respectively, while the transition remains steep and the ΔT_{hys} remains small (1–1.5 K).

The entropy changes as function of temperature, derived from magnetic isotherms through the Maxwell relations,^[5] are displayed in Figure 4. The absolute value of ΔS_m reaches $18 \text{ J kg}^{-1}\text{K}^{-1}$ around both 215 and 350 K, under a magnetic field change of 0–2 T. The peak values are rather stable (between 12.8 – $18.3 \text{ J kg}^{-1}\text{K}^{-1}$) throughout the whole temperature range from 220 to 380 K. These values are about 4 times greater than that of Gd (see the data included in Figure 4) for tuneable temperatures. Note that for the same effect more than twice the field change namely 0–5 T were required for $\text{MnFe}(\text{P}, \text{As})$ reported earlier.^[7] Thus, with the current materials much

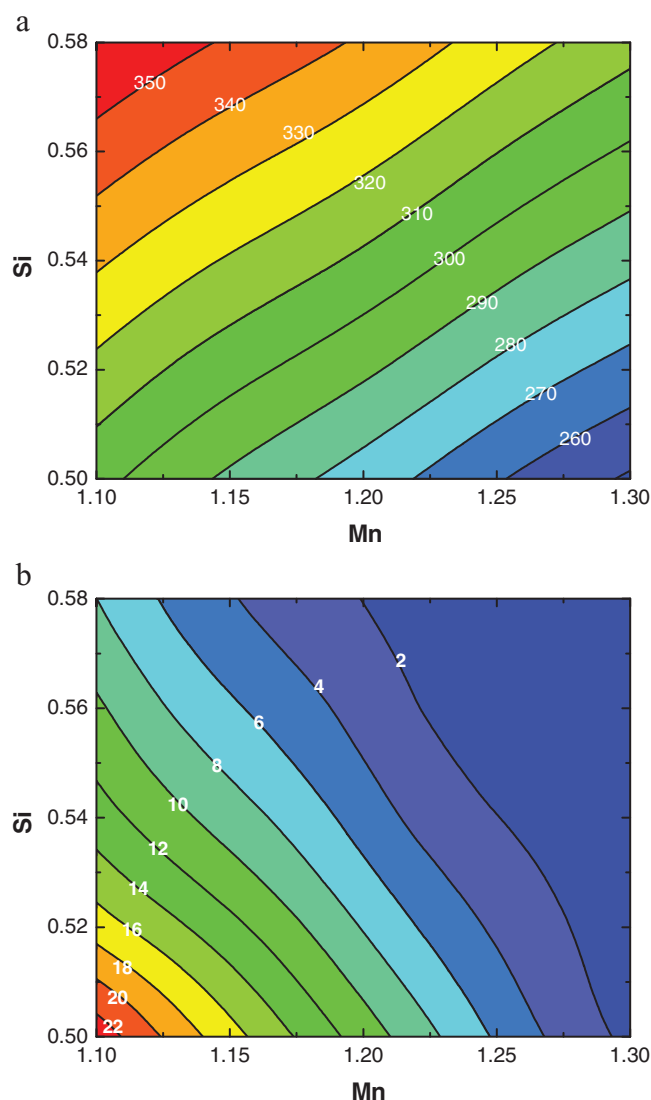


Figure 3. Partial phase diagram of the quaternary (MnFePSi) system (a) illustrating the composition dependence of the magnetic ordering temperature T_C (K) for $\text{Mn}_x\text{Fe}_{2-x}\text{P}_{1-y}\text{Si}_y$ compounds. (b) Composition dependence of the thermal hysteresis ΔT_{hys} (K) for $\text{Mn}_x\text{Fe}_{2-x}\text{P}_{1-y}\text{Si}_y$ compounds.

cheaper magnets may be employed in magnetocaloric refrigerators. Because the large effect is observed over a broad range of compositions, one can achieve an equally large MCE over a wide temperature interval by cascading several alloys with slightly different compositions in one active magnetic regenerator.^[13] Another important parameter to characterize the magnetocaloric effect is the adiabatic temperature change. From specific heat measurements in applied magnetic field we derive for $\text{Mn}_{1.24}\text{Fe}_{0.71}\text{P}_{0.46}\text{Si}_{0.54}$ an adiabatic ΔT of about 3 K at 320 K for a field change of 1 T. This result is very close to earlier results on $\text{Mn}_{1.1}\text{Fe}_{0.9}\text{P}_{0.47}\text{As}_{0.53}$ and Gd.^[4,5]

Because the effect is still large above the boiling point of water, the materials can be used in magnetocaloric generators (Figure 5) based on the Faraday induction law to transform abundant waste heat into electricity. This generator^[21]

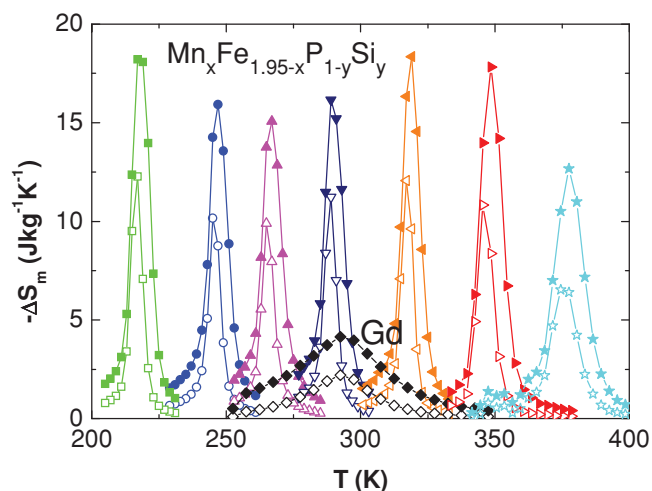


Figure 4. Isothermal magnetic entropy change under a field change of 0–1 T (open curves) and 0–2 T (solid curves) for some typical $\text{Mn}_x\text{Fe}_{1.95-x}\text{P}_{1-y}\text{Si}_y$ compounds with from left to right $x = 1.34, 1.32, 1.30, 1.28, 1.24, 0.66, 0.66$ and $y = 0.46, 0.48, 0.50, 0.52, 0.54, 0.34, 0.37$, respectively. The data of Gd metal under a field change of 0–1 T (open diamond) and 0–2 T (solid diamond) are included.

shall also contain a series of different materials as depicted in Figure 4 to utilize the full temperature span from room temperature up to the temperature of the heat source. As can be seen from Figure 4, to achieve a really constant magnetic entropy change, one will need to prepare many different materials with only slightly varying composition. Because the entropy change observed in a field change from 0–1 T is already rather large, this generator without moving parts can work efficiently. This is due to a magnetic-field-induced transition with a very small magnetic hysteresis, which occurs in these compounds at low fields, as displayed in Figure 6 for the compound $\text{Mn}_{1.30}\text{Fe}_{0.65}\text{P}_{0.50}\text{Si}_{0.50}$. The extremely small magnetic hysteresis

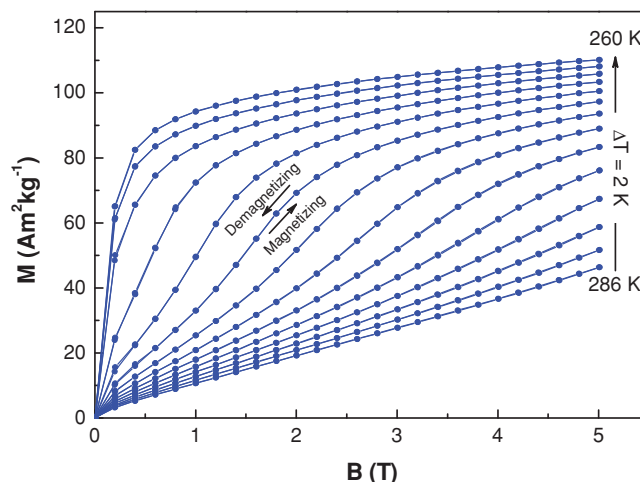


Figure 6. Magnetic isotherms of $\text{Mn}_{1.30}\text{Fe}_{0.65}\text{P}_{0.50}\text{Si}_{0.50}$ in the vicinity of the Curie temperature.

is in line with the observed small ΔT_{hys} , indicating a low energy barrier for nucleation of the FOMT.

3. Conclusions

Combining weak and strong magnetism in a single material opens the possibility to effectively couple spin and lattice degrees of freedom. In this way we simultaneously achieved giant magnetocaloric effect and a small thermal hysteresis in Mn–Fe–P–Si compounds of hexagonal Fe_2P -type structure by varying Mn:Fe, P:Si. We demonstrate that the working temperature can be controlled between 210 and 430 K by concurrently changing Mn:Fe and P:Si ratios. The combination of several alloys with slightly different compositions in one active magnetic regenerator will allow for efficient magnetic refrigeration with large temperature span. The fact that we use materials that are not only globally-abundant and non-toxic but can also be industrially-mass-produced via a simple powder-metallurgical method, makes Mn–Fe–P–Si compounds particularly attractive. The discovery of these high-performance low-cost magnetic refrigerants paves the way for commercialization of magnetic refrigeration and magnetocaloric power-conversion. Additionally, the insight into the importance of the coexistence of strong and weak magnetism enables us to search specifically for novel magnetocaloric materials.

4. Experimental Section

First principle electronic structure calculations were performed using the localized spherical wave method (LSW)^[22] using the scalar relativistic Hamiltonian. The local density exchange and correlation potential was used inside space filling and therefore overlapping spheres around the

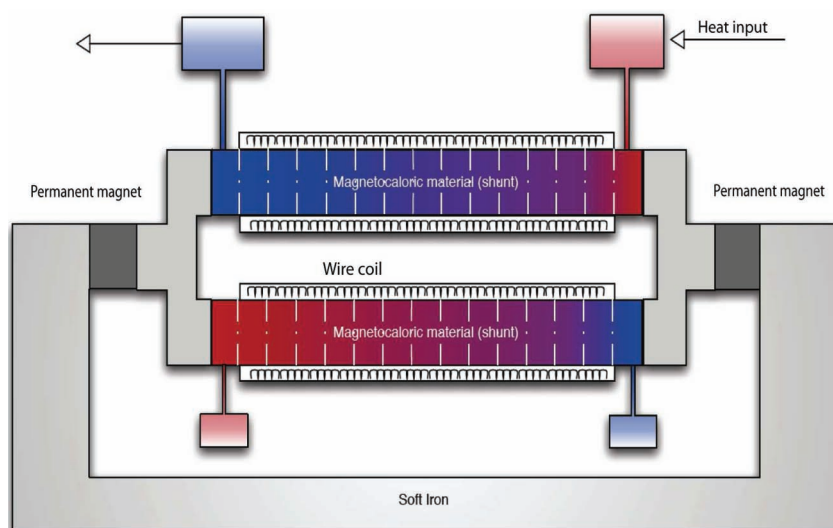


Figure 5. Concept of a magnetocaloric generator with an active shunt consisting of a series of materials with gradual changing Curie temperatures (adapted version of a design by Brillouin^[21]).

atomic constituents. Self-consistency was assumed when changes in the local partial charges were below 10^{-5} . The calculations simulating temperatures above the Curie-point employed supercells in the xy and z directions with antiferromagnetic (z) or ferrimagnetic (xy) starting magnetic configurations. Subtle mixing prevented falling back into the ground state. The visualization of the change in the electron density between the ferromagnetic and paramagnetic state used the VESTA^[23] program with VASP^[24] data as input.

Mn-Fe-P-Si compounds were prepared by ball milling and solid-state reactions from starting materials consisting of Mn (99.9%), Si (99.999%) chips, the binary compound Fe₂P (99.5%) and red-P (99.7%) powder. The powder obtained after milling for 10 hours was pressed into tablets. The tablets were sealed under Ar in quartz ampoules then sintered at 1373 K for 2 hours and annealed at 1123 K for 20 hours before slow cooling in the oven to room temperature. Powder X-ray diffraction patterns were obtained in a PANalytical X-pert Pro diffractometer with Cu K α radiation, secondary flat crystal monochromator and X'celerator RTMS Detector system. Magnetic measurements were carried out using the (Reciprocating Sample Option) RSO mode in a SQUID (Superconducting Quantum Interference Device) magnetometer (Quantum Design MPMS 5XL). The instrument's thermal lag of about 0.4 K with 1 K/min sweep rate around room temperature was obtained by measuring a Gd sample (3N Alfa Aesar) mounted in the same way as the other samples mentioned above. This thermal lag is not corrected for in the data displayed.

Acknowledgements

We thank Jack Voncken and Remco Addink for help with electron probe micro analysis and the artwork of Figure 5, respectively. This work is part of an Industrial Partnership Programme IPP I18 and Materials-specific Theory for Interface nad Nanoscience programme 88 of the 'Stichting voor Fundamenteel Onderzoek der Materie (FOM)' which is financially supported by the 'Nederlandse Organisatie voor Wetenschappelijk Onderzoek (NWO)'. The Industrial Partnership Programme is co-financed by BASF Future Business.

Received: May 13, 2011

Revised: July 17, 2011

Published online: September 9, 2011

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